



Zwicky's Missing Mass: Dark Matter versus Modified Gravity

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Abstract

Zwicky postulates the dark matter approach, based on the incompatibility between Newtonian gravity and the residual velocity measured in galaxy clusters. This approach to dark matter has been prioritized over other theoretical alternatives, although dark matter remains undetected directly. We discuss missing mass through large-scale modification of gravity, with the potential for an inverse Yukawa-type field (UYF), varying with comoving distance (null in the inner solar system, weakly attractive at interstellar distances, very attractive in galaxy clusters and repulsive on cosmic scales). The UYF is caused by the large-scale distribution of baryon mass, in the sense of Mach's principle. We obtain that the UYF resolved the incompatibility between the flatness of the Universe and the density of matter in Friedmann-Robertson-Walker cosmology, provides an explanation for dark energy as the cosmic acceleration caused by the large-scale distribution of matter. Also solves the Zwicky's missing mass because Virial's theorem includes an additional term, such that the virial mass is of the order of , five hundred times less than that calculated by Zwicky. X-ray emissions and Faber-Jackson and Tully-Fisher relationships are discussed.

Keywords: Zwicky Mass; Clusters of Galaxies; Dark Matter; Mach's Principle

Abbreviations

HST: Space Telescope; UYF: Yukawa-Like Field; BAO: Baryon Acoustic Oscillations; Λ FRW: Friedman-Robertson-Walker Cosmological Model.

Introduction

Zwicky's Missing Mass

In the 1930s, Fritz Zwicky reported a significant discrepancy between the visible mass of galaxy clusters and the mass necessary to explain their gravitational dynamics [1,2]. The application of the Virial theorem to the galaxies in the Coma cluster yielded a mass-to-light ratio of one

hundred times greater than that observed in the local group and in the solar stellar neighborhood [3]. While studying the Coma cluster, Zwicky realized that the galaxies were moving at such high speeds that, according to Newton's theory of gravitation, the visible mass of the cluster was not enough to hold them together. Furthermore, in the 1970s several studies were carried out on the dynamical motion of individual stars that suggest that the total mass of galaxies is about an order of magnitude greater than the mass of their visible parts [4], studies were also carried out on the rotation curves of galaxies showing incompatibility with the expected Keplerian motion in Newtonian gravity [5].

The problem of the incompatibility between astronomical observations and the formalism of gravitation turned the

dark matter conjecture into a paradigm after the observations of the anisotropies of the cosmic background radiation by the COBE satellite. Measurements of the angular size of the anisotropies showed that the Universe is asymptotically flat and, consequently, the observed density of matter should coincide with the critical density according to Einstein's field equations in the Friedmann-Robertson-Walker cosmology of the hot Big Bang. However the observed density of baryonic matter is an order of magnitude lower than expected.

All these astronomical observations, together with the measurements of gravitational lenses have in common their incompatibility with the formalism of gravitation on larger scales of the solar system [6]. After several decades of fruitless search for subatomic particles of exotic matter as a constituent of Dark Matter [7], and of new fundamental interactions (quintessence theories) the problem of missing mass has become a paradigm [8], similar to the ether paradigm, before the advent of special relativity in the 19th century, to the Phlogiston paradigm in the 18th century or even to the Ptolemy epicycles paradigm in the 16th century.

The lack of direct observation of the components of dark matter together with the manifestation of gravitational effects constitutes a paradox in itself. Zwicky's paradox is based on the assumption that Newton's Law is universally valid to describe the dynamics of the Universe, on astronomical scales, according to which the force of gravity is given by the inverse of the square of the distance and therefore the Virial theorem requires that the average kinetic energy be half the average potential energy. Let us remember that even General Relativity assumes Newton's Law of Gravitation as valid, considering that this is the limit to which the gravitational interaction tends in the weak field approximation.

On the other hand, the Inverse Square Law of distance assumes an infinite range for gravitational interaction, although Cosmology prescribes the finite radius of the universe, (Hubble radius) then an epistemological problem arises: How can an interaction have a range greater than the universe itself?

The infinite range of gravitation necessarily implies zero mass for the graviton, which contradicts the existence of detected gravitational waves. In the last two decades, the observation of gravitational waves emitted by the merger of a binary system was carried out through the LIGO-Virgo collaboration [9]. In these observations, the mass range of the graviton in the GW170104 event was determined to be on the order of 7.7×10^{-23} eV [9], and observation of binary black holes with space gravitational wave detectors was less constraining on the mass of the graviton less than 10^{-60} kg [10], it suggests that the range of the force of gravity is finite in contradiction to the Newtonian law of gravitation.

While it is true that the validity of inverse square law of Newton's gravity is verified with precisions greater than 10^{-8} for Eötvös-like experiments there is no empirical evidence of their validity beyond the solar system [11]; it is assumed true for estimating the mass of binary stars. When we use Newtonian gravity (the inverse square law of gravitation) to describe the dynamics of objects within the solar system, we use the two-body approximation and neglect the contributions of the other stars. This is justified, first of all, because the mass of the Sun is much greater than that of all the other components of the solar system and, consequently, the reduced two-body problem can be addressed perturbatively. And second, because the motion within the solar system is such that the comoving distance is negligibly small relative to the interstellar distance and therefore the gravitational contributions from the other stars in the galaxy are approximately the same at all points of the trajectory, that is, in a Gaussian sphere with a radius much smaller than the average interstellar distance (of the order of four light years). But these assumptions do not hold for comoving distance ranges on the order of kiloparsecs and megaparsecs (Figure 1). Within a galaxy cluster, the gravitational field caused by mass m_1 on mass m_2 (located at the edge of the yellow Gaussian) must also contain contributions from all other nearby galaxies, their total sum is not necessarily null, since galaxies within clusters do not have a homogeneously spherical distribution and comoving distances are not much smaller than the mean separation between galaxies.

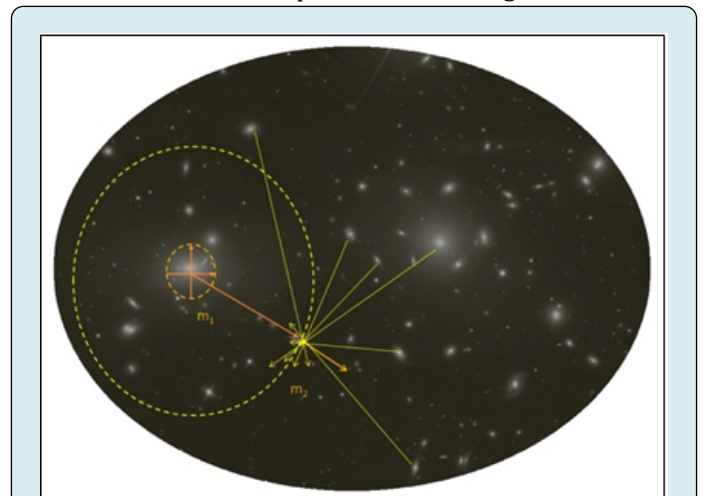


Figure 1: Application of March's Principle into the clusters of galaxies. The gravitational field of m_1 required the external contributions to the Gaussian sphere over m_2 . This complementary interaction is not prescribed by inverse square law of Newton's gravity.

So, local inertia is somehow linked to the large-scale distribution of matter in the Universe (March's Principle), then Newton's law of gravitation is insufficient to describe it. So the gravitational interaction between two stars or

between two galaxies would be the one prescribed by Newton's inverse square law plus an additional contribution from distant masses, which need not be the same for different points in space; since, on a large scale, the mass distribution is not spherical with its center in the local observation frame. Clearly, it is not possible to explicitly calculate that global contribution (due to the large-scale distribution of matter) to the gravitational force between two particles. Einstein tried it, through the cosmological term Λ , but it remained pending how to model his equivalent in stellar distances within a given galaxy, and within galaxy clusters.

Then, the general assumption is that all particles with mass are subject to the force of gravity through the inverse square law of gravitation, plus an additional term that varies with the comoving distance caused by the large-scale distribution of baryonic mass, in the sense of Mach's principle [12-15]. This complementary contribution would be zero at comoving distance ranges on the order of the solar system, weakly attractive at interstellar comoving distance ranges (kiloparsec), very attractive at intergalactic comoving distance ranges (Megaparsec) and repulsive at cosmic scales. It is proposed to address the Zwicky paradox of missing mass in galaxy clusters by modifying Newtonian gravity on a large scale with an inverse Yukawa-like term that varies with comoving distance [16].

To this end, a summary of modified large-scale gravity with a Yukawa-like inverse term, astrophysical implications

and consequences for FRW-cosmology are presented in Section 2. Next, we use the general derivation of the Clausius virial theorem and its application to data from the ESO nearby Abell cluster survey [17] (Section 3). The last section is a brief discussion on the linkage of the generalized Virial theorem (in the large-scale modification of Newtonian gravity) with the empirical relations between luminosity and residual velocity, in galaxies and galaxy clusters (Tully-Fischer and Faber-Jackson relations). Conclusions are given in the last section.

Methodology

Large-scale Modification of Newtonian Gravity

We assume that the net force of gravitation varies as the law of Inverse Square in scales in order of the interstellar distance, but it varies in a very different way when the comoving distance is about of the order of kiloparsec or more. In this sense, our argument is a large-scale modification of the gravity. The origin of this field (UYF) is the baryonic mass, like in the Newtonian gravity, and represents the contribution caused by the large-scale distribution of mass, in the sense of Mach's principle. This additional term was constructed phenomenologically to explain the astronomical observables. We call this term, due to its shape, as Inverse Yukawa-like field (UYF) [12].

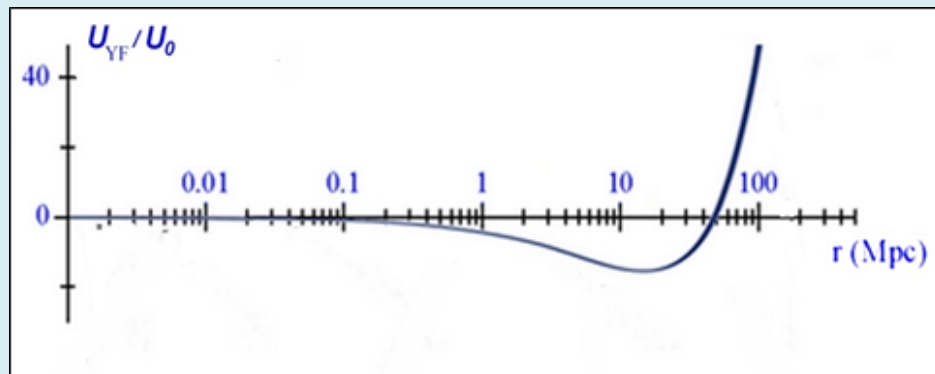


Figure 2: Modification at great scale of the gravity thought of the Inverse Yukawa-like field (UYF) in astronomical scale for different ranges of the comoving distance [12].

This potential per unit mass (in units of J/kg) as function of the comoving distance is: null in the inner solar system, weakly attractive in ranges of interstellar distances, very attractive in distance ranges comparable to the clusters of galaxies and repulsive to cosmic scales (Figure 1) [14]. The general expression is:

$$U_{YF}(s) \cong 2,5710^{45} \left[\frac{M}{M_{\odot}} \right] (s-1) e^{-0.05/s} \left(\frac{J}{kg} \right) \quad (1)$$

Where s is the dimensionless scale of the comoving distance $s \equiv r/r_0 = r/50\text{Mpc}$ and M is the mass of the object expressed in solar masses M_{\odot} . It's obtain beginning by the crude description of Hydrogen synthesis during the Matter-radiation decoupling in a primordial proto galaxies [15]. The coupling constants (as derived from previous report [15]) are: $r_0 \sim 50$ Mpc (the average distance between clusters of galaxies) and $\alpha \sim 2.5$ Mpc. We understand that an exact model would fit the precise values of the coupling constants without modifying the phenomenology.

The force per unit mass, complement to large-scale of the Newtonian gravitation (Figure 2), is:

$$F_{YF}(s) \cong -1.67 \cdot 10^{-21} \left[\frac{M}{M_{\odot}} \right] \frac{[s^2 + \alpha(s-1)]}{s^2} e^{-0.05/s} \left(\frac{N}{m} \right) \quad (2)$$

Notice that the maximum occurs in Abell radius, $r_m \cong 1.2$ Mpc as in the typical core of galaxies clusters.

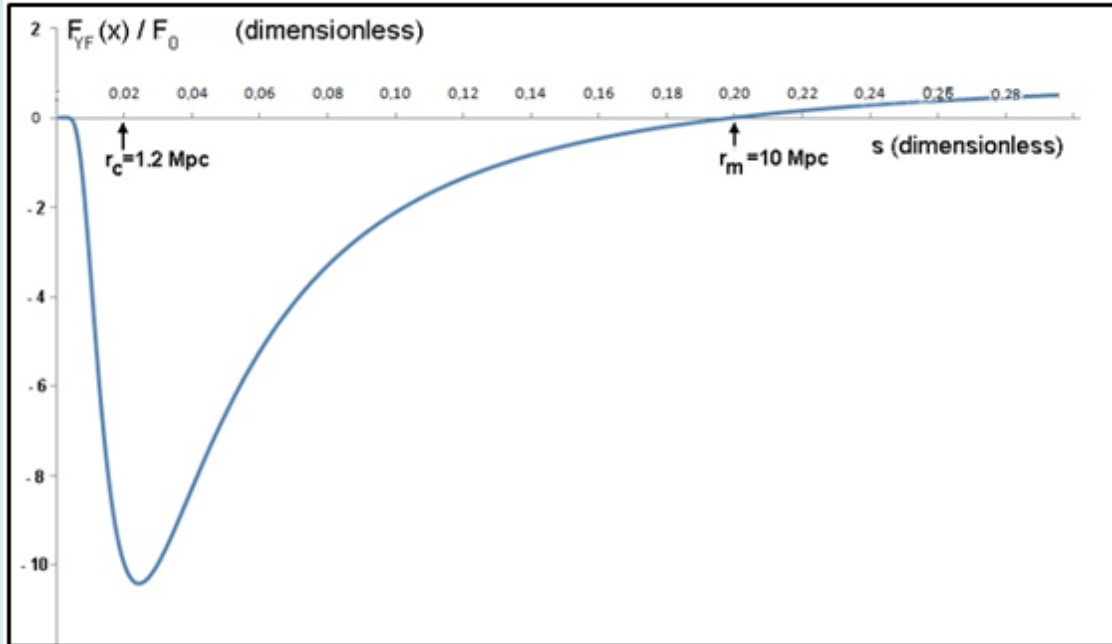


Figure 3: Force FYF in dimensionless scale of the comoving distances [15].

From Figures 2 & 3 it is clear that $U(r)$ gives a constant repulsive force per unit mass, at cosmological scales providing an asymptotic cosmic acceleration. This cosmic acceleration, on a large scale, remains constant as it is observed when taking the limit of s very large, for ranges of comoving distance much greater than $r_0 \sim 50$ Mpc.

For the average value of smooth transition to strong agglutination in galaxy's distribution ($r_c \sim 10$ Mpc) the FYF is null, it suggests that the range of the force of gravity is finite; and the graviton rest mass is $m_g^0 \cong 10^{-29} eVc^{-2}$ [14], in according to the results of gravitational waves [10]. The large scale structures with characteristic dimensions much greater than 10 Mpc; e.g. Sloan Great Wall and Voids; do not show symmetric axial distribution that would be expected if gravitation had infinite range. Neither have spherical distribution the hot gas in the super clusters of galaxies found by means of the Suyaev-Zel' dovich. In large-scale structures, the gravitational bond between the galaxies, maybe by a sequential chain of gravitational attractions between their neighboring components, but not necessarily by a common center. Assuming an infinite range for gravity, would imply among other things, to imagine colossal masses for the attractor center in the super clusters of galaxies, which are unobservable (Hypermassive Black Hole).

Cosmological Implications

Let us now consider a usual Friedman-Robertson-Walker cosmological model (Λ FRW), with homogeneous and isotropic FRW-metric together energy-momentum tensor for a perfect fluid, then we obtain the usual Friedmann equations [12-14] with cosmological constant L and null curvature ($k=0$):

$$\left(\frac{\dot{R}(t)}{R(t)} \right)^2 = \frac{8\pi G}{3} \rho + \frac{\Lambda c^2}{3} \quad (3)$$

$$\frac{2\ddot{R}(t)}{R(t)} + \left(\frac{\dot{R}(t)}{R(t)} \right)^2 = -\frac{8\pi G}{c^2} P + \Lambda c^2 \quad (4)$$

Where $R(t)$ is the Riemann curvature tensor, P and ρ denote the pressure and density respectively. As usual, c denotes the speed constant of light and G the acceleration constant of gravity.

Now, we assumed Λ as a cosmic variable respect to the comoving distance. Note that the covariance is guaranteed because at cosmological scales (ranges of the comoving distance: $r > 50$ Mpc) the F_{YF} is constant (Figure 3). Then in these cosmological scales, the galaxies are described

as dust particles through the impulse energy tensor for perfect fluid. Thus the Dark Energy can be thought of as a “cosmic force” in the sense of the Mach Principle, caused by ordinary matter, through the Lambda Λ cosmological term.

$$\Lambda \equiv \Lambda_0 F_{YF}(r) = -\frac{3H_0}{c^3} \frac{d}{dr} U_{YF}(r) \quad (5)$$

By other hand, when $r \rightarrow r_m$, i.e. $r_m \cong 1.2$ Mpc (Figure 3), and using (1) into (5), we obtain $L(r)$ in the intergalactic scale, as:

$$\Lambda(r_m) = \Lambda_0 F_{YF}(r) \Big|_{r \rightarrow r_m} \cong 10.55 (4\pi G \text{ kg m}^{-2}) \frac{3H_0}{c^3} \quad (6)$$

With H_0 is the Hubble constant in km/s Mpc⁻¹. Remember that $s \equiv r/r_0 = r/50\text{Mpc}$. Now, the Friedmann equations are [14];

$$\left(\frac{\dot{R}(t)}{R(t)} \right)^2 = \frac{8\pi G}{3} \rho + \frac{\Lambda(r_m)c^2}{3} \quad (7)$$

$$\frac{2\ddot{R}(t)}{R(t)} + \left(\frac{\dot{R}(t)}{R(t)} \right)^2 = -\frac{8\pi G}{c^2} P + \Lambda(r_m)c^2 \quad (8)$$

The Friedmann equations are:

$$\frac{kc^2}{R^2(t)} = H_0^2 [\Omega_b(1 + \Omega_{YF}) + \Omega_\Lambda - 1] \quad (9)$$

$$q_0 = \frac{\Omega_b}{2} (1 + \Omega_{YF}) (1 + \frac{3P}{c^2 \rho}) - \Omega_\Lambda \quad (10)$$

Where and we used the standards notation for the matter density Ω_m , baryonic density Ω_b , and deceleration parameters Ω_Λ .

The remarkable result in (9) is that: if $k=0$ no need to the dark matter. I.e. using $\Omega_b \approx 0.0223$ and $\Omega_\Lambda \approx 0.6911$ [7,15] we obtain $\Omega_m \approx 0.255$ and $\Omega_m + \Omega_\Lambda = 0.255 + 0.6911 \approx 1$.

Now consider the photons emitted from a remote galaxy with recession velocity v , and their observation in the reference local frame. We should evaluate (2) in $r \gg 50$ Mpc, with initial condition $v(0)=0$. We obtain de Hubble-Lemaître's Law [13]:

$$v = \int adt = \int \left(\lim_{r \rightarrow \infty} F_{YF}(r) \right) \frac{dr}{c} = \left(\frac{4\pi G}{c} \ell \right) r = H_0 r \quad (11)$$

here $\ell \equiv 1 \text{ kg m}^{-2}$ is a dimensional parameter.

Notice that the value of H_0 is the theoretical upper limit, evaluate for most distant objects ($r \gg 50$ Mpc). Planck

collaborations obtain indirect measurement $67.15 \text{ km s}^{-1} \text{ Mpc}^{-1}$; Riess had found that $74.22 \text{ km s}^{-1} \text{ Mpc}^{-1}$ in Large Magellanic Cloud and other recent direct measurements which Space Telescope (HST), are 75.8 and $78.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$ [18,19].

In the cosmological range of the comoving distance, $r^{\otimes} r_c$, i.e. when $s^{\otimes} 0.2$ (Figure 3), the L parameter using (1) into (5), is

$$\Lambda(r_c) = \Lambda_0 F_{YF}(r) \Big|_{r \rightarrow r_c} \cong 0.623 \frac{3H_0^2}{c^2} \quad (12)$$

And the Friedmann equations are

$$\left(\frac{\dot{R}(t)}{R(t)} \right)^2 = \frac{8\pi G}{3} \rho + \frac{\Lambda(r_c)c^2}{3} \quad (13)$$

$$\frac{2\ddot{R}(t)}{R(t)} + \left(\frac{\dot{R}(t)}{R(t)} \right)^2 = -\frac{8\pi G}{c^2} P + \Lambda(r_c)c^2 \quad (14)$$

Thus, the dark energy would be interpreted as the cosmic acceleration in local frameworks, caused by the large scale distribution of the ordinary baryonic matter. Thus, the cosmological density parameter is now $\text{WL@ } 0.623h^{-1}$. The upper limit for Hubble parameter ($h=0.863$), we obtain $\Omega_\Lambda \approx 0.72$ in good agreement with the measurements of SNIa [14].

Other important astrophysical implications have been presented in previous reports: we discussed how the UYF affects the calculation of the age of the universe, the rotation curves of galaxies, the Angular Diameter Distance Distributions, and the length and Jeans's mass [14]; also the inclusion of the UYF in Arp Controversy, the Gravitational redshift. Gravitational Lensing, baryon acoustic oscillations (BAO) and Cosmic Microwave Background anisotropies (CMB) and Pioneer Anomaly were presented in [15].

Result

Clausius's Virial Theorem

Let us consider the Clausius's Virial expression $G \equiv \sum_i \vec{p}_i \cdot \vec{r}_i$ when deriving with respect to time, and then averaging with respect to a complete period (τ), we obtain the well-known virialized expression between kinetic energies and power; that in the case of a particle subjected only to the Newtonian gravitational potential is:

$$\begin{aligned} \frac{1}{\tau} \int_0^\tau \frac{dG}{dt} dt' &= \frac{1}{\tau} \int_0^\tau \sum_i -\vec{\nabla} U_i \cdot \vec{r}_i dt' + \frac{1}{\tau} \int_0^\tau \sum_i \frac{p_i^2}{2m_i} dt' \\ \frac{G(\tau) - G(0)}{\tau} &= -\frac{1}{\tau} \int_0^\tau \sum_i \frac{GMm_i}{r_i} dt' + \frac{1}{\tau} \int_0^\tau \sum_i T_i dt' \\ 0 &= -\langle U \rangle + 2\langle T \rangle \end{aligned} \quad (16)$$

where T is the kinetic energy. The virial theorem is often used to calculate the mass-to-light ratio in the Zwicky missing mass problem. Using the projected virial radius (R_{pv}) and projected radial velocity (σ_p) [17], then

$$M = \frac{3\pi}{2} \frac{\sigma_p^2}{G} R_{pv} \quad (17)$$

If now the particles are subjected to a gravitational potential that have an additional long-range term, we have using (1) and (2) that:

$$\begin{aligned} 0 &= -\frac{GMm}{r} - 4\pi G M m \ell (r - r_0) e^{-\alpha/r} + 2T + \frac{4\pi G M m \ell}{\tau} \int_0^r \left(\alpha - \frac{\alpha r_0}{r} \right) e^{-\alpha/r} dt \\ 0 &= -\frac{GMm}{r} - U_{VF} + 2T + \frac{4\pi G M m \ell}{\tau} \int_0^r \left(\alpha - \frac{\alpha r_0}{r} \right) e^{-\alpha/r} dt \end{aligned} \quad (18)$$

Remember that $s \equiv r/r_0 = r/50Mpc$, $\alpha_0 \equiv \alpha/r_0$ and $\ell \equiv 1 \text{ m}^{-2}$ is a dimensional parameter.

Thus, the energy balances results:

$$v^2 = GM \left[\frac{1}{r_0 s} - 4\pi \ell r_0 \left(e^{-\alpha/r} (s-1) + \alpha_0 \right) \right] \quad (19)$$

Then the generalized expression of the Virial Theorem in a cluster of galaxies of radius R is:

$$v^2 = GM R^{-1} \left[1 - 4\pi \ell r_0^2 s \left(e^{-\alpha_0/s} (s-1) + \alpha_0 \right) \right] \quad (20)$$

A preliminary calculation of virial mass is carried out for some nearby bright galaxy clusters in the Abell Catalogue, together with data reported in optical [17]. Then (20) can be writing as:

$$M = \frac{3\pi}{2} \frac{\sigma_p^2}{G} R_{pv} \left[1 - 4\pi \ell R_{pv} \left(e^{-2.5/R_{pv}} (R_{pv} - 50) + 2.5 \right) \right]^{-1} \quad (21)$$

Where σ_p is the radial component speed for spherical systems, and R_{pv} is the projected virial radius.

Observation data for the bright Abell galaxy clusters are shown in Table 1. The galaxy clusters are listed in the first column, followed by the radii, redshift, radial component speed and the projected virial radius in the second to fifth columns respectively. In the penultimate column, we have shown the calculation of mass according to Newtonian gravitation. And in the last column we show the modified gravity mass. Note that the estimated mass is now, on average, 500 times less than originally calculated [17]; thus resolving the Zwicky paradox without invoking non-baryonic dark matter.

Cluster	R_{\max} (Mpc) [17]	Z [17]	σ_p km/s [17]	$R_{pv} \text{ h}^{-1}$ (Mpc) [17]	M_{vir} (17) ($10^{14} M_{\odot}$)	$M_{\text{vir}}^{(VF)}$ (21) ($10^{11} M_{\odot}$)
A13	1.17	0.0972	515	1.03	2.99	4.45
A16	1.94	0.0467	759	1.28	8.07	9.22
A26	1.2	0.0693	179	0.27	0.1	0.584
A37	1.1	0.0389	508	0.76	2.15	4.55
A85	1.75	0.056	969	1.31	13.5	14.9
A87	0.74	0.0558	224	0.55	0.3	0.91
A151	1.56	0.0537	714	0.88	4.91	8.8
A168	1.1	0.0469	118	0.25	0.04	0.25
A194	1.97	0.0184	341	0.55	0.7	2.1
A229	1.31	0.1137	506	1	2.8	4.32
A256	1.8	0.0885	545	0.56	1.82	5.37
A262	3.95	0.0169	525	0.87	2.63	4.77
A295	1.4	0.0427	359	0.64	0.9	2.31
A367	1	0.0882	394	0.68	1.16	2.77
A399	1.56	0.0718	1116	1.19	16.2	20.3
A400	1.22	0.0237	599	0.7	2.75	6.38
A401	2	0.0737	1152	1.1	16	22
A420	1.05	0.086	360	0.67	0.95	2.31
A426	2.7	0.0178	1026	1.35	15.6	16.6

A458	0.88	0.1057	736	1.1	6.52	8.97
A514	1.77	0.0714	882	1.05	8.94	13
A524	0.89	0.0797	250	0.37	0.25	1.14
A539	2.47	0.0284	629	0.68	2.95	7.06
A548NE	0.88	0.0397	571	0.76	2.71	5.75
A548SW	1.36	0.0439	583	0.77	2.87	5.98
A569	3.03	0.0201	327	0.46	0.54	1.95
A576	1.74	0.0383	914	1.22	11.2	13.5
A754	2.6	0.0539	662	0.98	4.7	7.43
A957	5.17	0.0416	254	0.3	0.21	1.18
A1016	1.44	0.0319	273	0.4	0.33	1.36
A1060	2.06	0.0126	610	0.69	2.81	6.63
A1069	1.2	0.0662	360	0.61	0.87	2.33
A1142	2.4	0.035	486	0.93	2.41	4.04
A1146	1.72	0.1422	929	0.84	7.94	15
A1185	3.97	0.03	536	0.38	1.2	5.24
A1314	1.17	0.0329	277	0.47	0.4	1.4
A1367	3.08	0.0221	570	0.73	2.6	5.75
A1631	1.95	0.0464	702	1.13	6.1	8.11
A1651	1.23	0.0863	685	0.5	2.57	8.52
A1656 Coma	3.9	0.0233	821	1.08	7.97	11.2
A1689 Virgo	2.19	0.0038	632	1.05	4.59	6.68
A1736	1.82	0.0347	415	0.92	1.73	2.95
A1795	1.81	0.0632	834	0.91	6.93	11.9
A1809	1.65	0.079	765	0.79	5.06	10.3
A1983	1.7	0.0452	494	0.79	2.11	4.28
A1991	0.81	0.0593	631	0.48	2.09	7.24
A2029	1.57	0.0766	1164	0.48	7.12	24.6
A2040	1.85	0.0454	458	0.92	2.11	3.6
A2048	1.19	0.0972	664	0.84	4.06	7.67
A2052	1.15	0.0346	207	0.47	0.21	0.78
A2063	3.87	0.035	667	0.66	3.22	7.95
A2079	2.76	0.0662	670	1.12	5.51	7.41
A2092	1.64	0.0673	536	0.45	1.42	5.23
A2107	1.02	0.0415	622	0.71	3.01	6.87
A2124	1.22	0.0661	878	0.88	7.43	13.3
A2199	3.29	0.0314	801	0.96	6.74	10.9
A2256	1.19	0.0589	1348	1.32	26.3	28.8
A2319	1.04	0.0553	1545	1.83	47.8	34.8
A2353	1.9	0.1213	597	0.61	2.38	6.41
A2362	0.81	0.0616	331	0.57	0.68	1.98

A2401	0.63	0.0581	395	0.69	1.18	2.78
A2426	1.35	0.0886	332	0.49	0.59	2
A2500	0.96	0.0904	477	0.58	1.45	4.1
A2538	0.74	0.0858	326	0.68	0.79	1.9
A2554	0.68	0.1118	840	0.82	6.34	12.3
A2569	1.16	0.0816	491	0.88	2.32	4.16
A2589	0.59	0.0423	470	0.33	0.8	4.03
A2593	1.41	0.0424	698	0.43	2.29	8.88
A2634	0.86	0.0316	700	1.03	5.53	8.23
A2666	0.55	0.028	383	0.29	0.47	2.68
A2670	2.22	0.0767	852	0.99	7.87	12.3
A2715	1.22	0.1145	463	0.85	2	3.72
A2717	1.16	0.0498	541	0.8	2.56	5.13
A2721	1.6	0.1152	805	1.01	7.17	10.9
A2734	1.77	0.0625	628	0.98	4.23	6.69
A2755	1.23	0.0957	768	1.21	7.81	9.56
A2798	2.21	0.1129	711	0.41	2.27	9.21
A2799	0.73	0.064	422	0.75	1.46	3.14
A2800	0.96	0.0643	404	0.78	1.39	2.87
A2819	2.01	0.0876	282	0.64	0.56	1.43
A2854	0.77	0.0619	130	0.25	0.05	0.31
A2877	1.08	0.0248	887	0.59	5.08	14.2
A2911	1.06	0.0816	547	0.72	2.36	5.31
A3093	0.9	0.0836	440	0.85	1.8	3.36
A3094	1.71	0.0677	653	1.26	5.88	6.85
A3111	1.08	0.0774	159	0.24	0.06	0.46
A3112	1.2	0.0735	86	0.13	0.01	0.13
A3122	1.76	0.0648	775	1.1	7.23	9.95
A3126	0.58	0.0862	1053	0.67	8.13	19.8
A3128	2.42	0.0605	789	1.28	8.73	9.96
A3142	1.13	0.1036	737	1.02	6.07	9.14
A3151	0.56	0.0662	237	0.4	0.25	1.02
A3158	1.67	0.0597	976	1.06	11	15.9
A3182 Eridani	1.94	0.0058	264	0.72	0.55	1.24
A3194	1.25	0.0977	805	1.2	8.52	10.5
A3223	1.74	0.0603	647	0.87	3.99	7.24
A3266	1.34	0.0599	1107	1.91	25.6	17.6
A3202	0.74	0.0708	250	0.68	0.47	1.11
A3334	0.54	0.097	696	0.85	4.51	8.41
A3341	0.93	0.0389	351	0.58	0.78	2.22

A3354	1.68	0.0586	358	0.72	1.01	2.27
A3360	0.49	0.0848	835	1.46	11.1	10.8
A3376	2.29	0.0465	688	0.89	4.61	8.16
A3381	1.6	0.0382	293	0.45	0.42	1.56
A3389	0.52	0.0272	595	0.74	2.87	6.26
A3391	0.87	0.0553	663	0.92	4.43	7.54
A3395	1.14	0.0506	852	0.76	6.04	12.8
A3526	4.56	0.0108	447	0.78	1.71	3.51
A3528N	1.71	0.0547	461	0.56	1.3	3.84
A3532	0.66	0.0559	738	1.19	7.1	8.86
A3556	1.04	0.0477	642	1.02	4.6	6.93
A3558 Shapley	1.99	0.048	977	1.25	13.1	15.4
A3559	1.75	0.0469	456	0.74	1.68	3.68
A3562	2.15	0.0478	736	1.22	7.24	8.77
A3571	0.98	0.0395	1045	0.73	8.73	19.3
A3574	1.26	0.0158	491	0.8	2.11	4.22
A3651	1.8	0.061	626	1.11	4.76	6.48
A3667	2.22	0.0566	971	1.55	16	14.4
A3693	1.06	0.0921	478	0.86	2.15	3.96
A3695	1.97	0.0903	779	1.3	8.64	9.67
A3703	1.04	0.0743	472	0.96	2.34	3.79
A3705	1.08	0.0906	877	1.07	9.01	12.8
A3716N	1.42	0.0493	466	0.78	1.85	3.82
A3716S	1.91	0.0458	803	0.85	6	11.2
A3733	0.72	0.0399	610	0.92	3.75	6.38
A3806	2.17	0.0762	502	1.1	3.04	4.17
A3809	1.48	0.0631	478	0.9	2.25	3.93
A3822	1.86	0.0769	810	1.49	10.7	10.1
A3825	1.66	0.076	699	1.2	6.42	7.94
A3879	1.74	0.0679	398	0.5	0.87	2.88
A3880	0.79	0.0588	827	0.92	6.89	11.7
A3921	1.33	0.0944	490	0.93	2.45	4.11
A4008	0.69	0.0558	427	0.66	1.32	3.26
A4010	1.44	0.0966	625	0.88	3.76	6.75
A4038	0.42	0.0285	413	0.43	0.8	3.11
A4067	1.45	0.0998	499	0.83	2.26	4.34
Average	1.63	0.0606	606	0.83	4.73	7.3

Table 1: Virial Mass in Abell bright clusters of galaxies.

Discussion

New Cosmological Paradigm

The total gas mass in clusters is generally larger than the total mass of the luminous parts of the galaxies $M_{\text{gas}} \sim 10^{13}\text{-}10^{14} M_{\odot}$ [20]. The increase in the effective potential, through the UYF term in the Virial equation (20), decreases the virial mass of the galaxy clusters, necessary to maintain hydrostatic balance, consequently the gas fraction and the baryon density increase, without the need of assuming halos of exotic dark matter (non-baryonic dark matter), as in the Table 1. Also, we can see (Figure 4) that the correction in the projected radial velocity is the order of 10^{-3} for the comoving distance ranges in galaxy clusters.

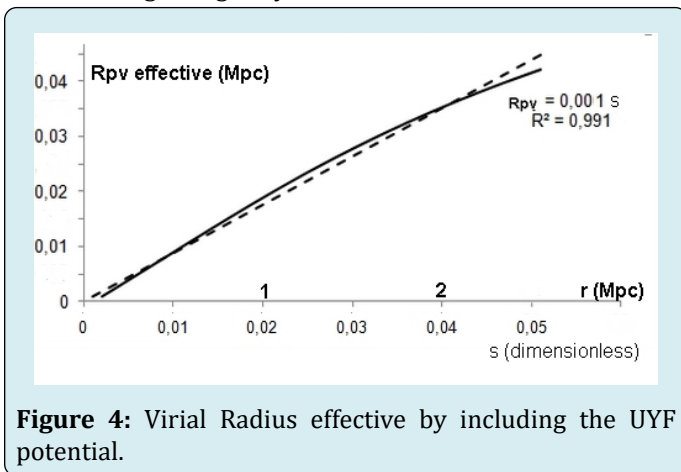


Figure 4: Virial Radius effective by including the UYF potential.

Yet, in the atmosphere of galaxy clusters, the hot gas emitting X-rays. Given the old of the galaxies clusters, and considering that global eruptive events in most galaxies are sporadic, we can assume hydrostatic equilibrium, therefore the energy balance prescribes that the temperature of the gas follows the variation of velocities, analogously to the distribution of matter, then:

$$k_B T_{\text{eff}} \cong \mu m_p \sigma_p^2 \quad (22)$$

where μ is the metallicity of the order 0.6 times that of the solar, m_p is the mass of the proton, and k_B is the Boltzmann constant. The Virial theorem (20) implies that

$$T_{\text{eff}} \cong \left(\mu m_p k_B^{-1} G r_0^{-1} \right) M s^{-1} \left[1 - 4\pi \ell r_0^2 s \left(e^{-\alpha_0/s} (s-1) + \alpha_0 \right) \right] \quad (23)$$

The plasma temperature turns out to be non-uniform (Figure 5), but beyond the core of the galaxy cluster the isothermal hypothesis is valid. On the other hand, the mass of the gas, deduced from X-ray observations and the Sunyaev-Zel'dovich effect, gives masses of the order of $10^{13} - 10^{14} M_{\odot}$, and $T_{\text{eff}}^{(X\text{-ray})} \approx 210^7 - 10^8 K$ [20]. So consideration of the Virial Theorem in a modified gravity scenario predicts that the gas fraction is much larger than the mass of the stars and,

therefore, the Baryon density parameter increases according to the Friedmann's equation (9) without to invoke exotic dark matter.

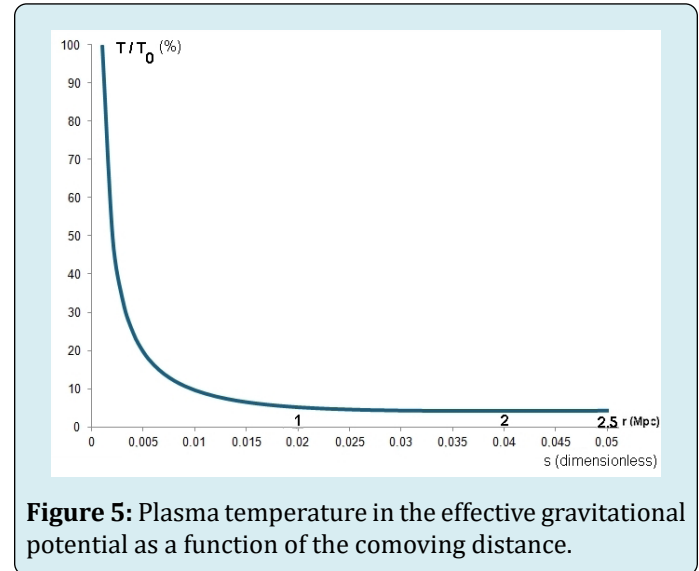


Figure 5: Plasma temperature in the effective gravitational potential as a function of the comoving distance.

On the other hand, the Faber-Jackson relation [21] is the observed relationship between the luminosity (L) and the random speed (σ) of stars near to center elliptical galaxy: $L \propto \sigma^4$. This empirical relationship can be deduced from the virial theorem for Newtonian gravitation. In effect using (16):

$$\frac{3}{5} \frac{GM^2}{R} = M\sigma^2 \quad (24)$$

And the assumption that $L \propto M$, and $L = 4\pi R^2 B$, with B superficial briskness, then:

$$L = 4\pi \left(\frac{3}{5} \frac{GL}{\sigma^2} \right)^2 B \rightarrow L \propto \sigma^4 \quad (25)$$

Let us now consider the large-scale modification of Newtonian gravity, using the generalized Virial theorem (20) into (26) then:

$$L = 4\pi \left(\frac{3}{5} \frac{GL}{\sigma^2} \right)^2 \left[1 - 4\pi \ell r_0^2 s \left(e^{-\alpha_0/s} (s-1) + \alpha_0 \right) \right]^2 B \rightarrow L \propto \sigma^4 \quad (26)$$

Also in spiral galaxies there is an empirical relationship between the mass or intrinsic luminosity of a spiral galaxy and its asymptotic rotation velocity or emission line width; known as the Tully-Fisher relation [22]. More precisely, the baryonic Tully-Fisher relation says that the baryonic mass is proportional to velocity to the power of roughly 3.5-4, [23,24]. Therefore, the Tully-Fisher relationship suggests that the luminous mass in stars and gas accounts for almost all of the associated baryons in each individual galaxy and its halo. Therefore this is a strong argument against dark

matter in these systems [25]. The generalization of the Virial Theorem in an scenarios of modified gravity allows know the basic physics in Tully–Fisher and Faber-Jackson relations (26) without invoke the paradigm of dark matter.

Finally, a note is necessary on the validity of Mach's Principle in gravitation (Figure 1). The general solution to the gravitation Poisson equation, under spherical symmetry, depends on the mass distribution outside of r :

$$U(r) = -4\pi G \left[\frac{1}{r} \int_0^r \rho(r') r'^2 dr' + \int_r^\infty \rho(r') r' dr' \right] \quad (27)$$

Newton's first theorem: a body that is inside a spherical shell of matter experiences no net gravitational force from that shell. The equivalent in general relativity is called Birkhoff theorem. In the Newtonian approach to gravity it is easily understood that the second term in (27) cancels due to the fact that solid angles, extending from a point inside a sphere in opposite directions, have areas on the sphere that vary as r^2 . While the angles sustained by the gravitational forces vary as r^2 , so that the gravitational forces of the two opposite areas exactly cancel. But, when the gravitational force is different from Newton's inverse square law of gravity, the two opposite areas no longer change as r^2 and, consequently, it is not true that the second term in (15) is null. Therefore, Birkhoff theorem could not be applied.

Summary and Conclusions

Almost a century ago, Zwicky introduced the dark matter approach, based on the incompatibility between Newtonian gravity and the residual velocity measured in galaxy clusters. The incompatibility between the observed matter density with the flatness of the Universe in FRW-cosmology; and non-Keplerian rotation-speed curves in galaxies and galaxy clusters, show that our formulation of gravity is far from complete. However, the dark matter approach has been prioritized over other theoretical alternatives (quiescence, bigravity, modified gravity, etc.), turning dark matter into an uncritically assumed paradigm, despite the lack of detection of exotic matter, in terrestrial laboratories and astronomical observations, during of the last thirty years.

A plausible alternative is the large-scale modification of gravity, through a potential as an inverse Yukawa-like field (UYF-field), which varies with comoving distance (null in the inner solar system, weakly attractive at interstellar distances, very attractive in galaxy clusters and repulsive on cosmic scales). The UYF-field is caused by the large-scale distribution of baryon mass, in the sense of Mach's principle, and can explain astronomical observables, particularly

- Resolved the incompatibility between the flatness of the Universe and the density of matter in the FRW-Cosmology without dark matter (9).

- Provides a explanation for dark energy (14), The dark energy would the cosmic acceleration caused by the large scale distribution of the ordinary baryonic matter, as prescribed the Mach's principle.
- It allows us to theoretically deduce the Hubble-Lemaître's law (11).
- The Virial theorem and Kepler's third law now including additional terms (18) that solve the Zwicky's paradox, so the "missing mass" is reinterpreted as the potential energy associated with the UYF-field (21).
- The virial mass in the galaxy cluster is the order of $10^{11} M_\odot$. This is five hundred times less than that calculated by Zwicky (Table 1).
- It is compatible with the empirical relations of Faber-Jackson and Tully-Fisher (25).

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